Reading **Required:** Witkin, Particle System Dynamics, SIGGRAPH '01 course notes on Physically Based Modeling. (online handout) • Witkin and Baraff, Differential Equation Basics, **Particle Systems** SIGGRAPH '01 course notes on Physically Based Modeling. (online handout) Optional **Brian Curless** Hockney and Eastwood. Computer simulation **CSEP 557** using particles. Adam Hilger, New York, 1988. Spring 2019 • Gavin Miller. "The motion dynamics of snakes and worms." Computer Graphics 22:169-178, 1988. 2 1 Particle in a flow field What are particle systems? We begin with a single particle with: A particle system is a collection of point masses that • Position, $\mathbf{x} = \begin{bmatrix} x \\ y \end{bmatrix}$ obeys some physical laws (e.g, gravity, heat convection, spring behaviors, ...). • Velocity, $\mathbf{v} \equiv \dot{\mathbf{x}} = \frac{d\mathbf{x}}{dt} = \begin{bmatrix} dx / dt \\ dy / dt \end{bmatrix}$ Particle systems can be used to simulate all sorts of physical phenomena: Suppose the velocity is actually dictated by a driving Sum function, a vector flow field, g: van $\dot{\mathbf{x}} = \mathbf{g}(\mathbf{x},t)$ Fire volnutric light Eloth Smoke dust jello astronnica V х bees If a particle starts at some point in that flow field, how should it move? 3

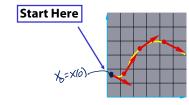
Diff eqs and integral curves

The equation

 $\dot{\mathbf{x}} = \mathbf{g}(\mathbf{x}, t)$

is actually a first order differential equation.

We can solve for x through time by starting at an initial point and stepping along the vector field:



This is called an **initial value problem** and the solution is called an **integral curve**.

Particle in a force field

Now consider a particle in a force field **f**.

In this case, the particle has:

Mass, m

• Acceleration, $\mathbf{a} = \mathbf{\ddot{x}} = \mathbf{\dot{v}} = \frac{d\mathbf{v}}{dt} = \frac{d^2\mathbf{x}}{dt^2}$

The particle obeys Newton's law:

 $\mathbf{f} = m\mathbf{a} = m\ddot{\mathbf{x}}$

So, given a force, we can solve for the acceleration:

 $\ddot{\mathbf{x}} = \frac{\mathbf{f}}{m}$

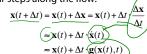
The force field \mathbf{f} can in general depend on the position and velocity of the particle as well as time.

Thus, with some rearrangement, we end up with:

$$\ddot{\mathbf{x}} = \frac{\mathbf{f}(\mathbf{x}, \dot{\mathbf{x}}, t)}{m}$$

Euler's method

One simple approach is to choose a time step, Δt , and take linear steps along the flow:



Writing as a time iteration:

$$\mathbf{x}_{\nabla}^{i+1} = \mathbf{x}^{i} + \Delta t \cdot \mathbf{g}^{i} \quad \text{with} \quad \left(\mathbf{g}^{i}\right) = \mathbf{g}(\mathbf{x}), t = i\Delta t)$$

This approach is called **Euler's method** and looks like:



Properties:

- Simplest numerical method
- Bigger steps, bigger errors. Error ~ $O(\Delta t^2)$.

Need to take pretty small steps, so not very efficient. Better (more complicated) methods exist, e.g., adaptive timesteps, Runge-Kutta, and implicit integration.

Second order equations

This equation:

$$\ddot{\mathbf{x}} = \frac{\mathbf{f}(\mathbf{x}, \dot{\mathbf{x}}, t)}{m}$$

is a second order differential equation.

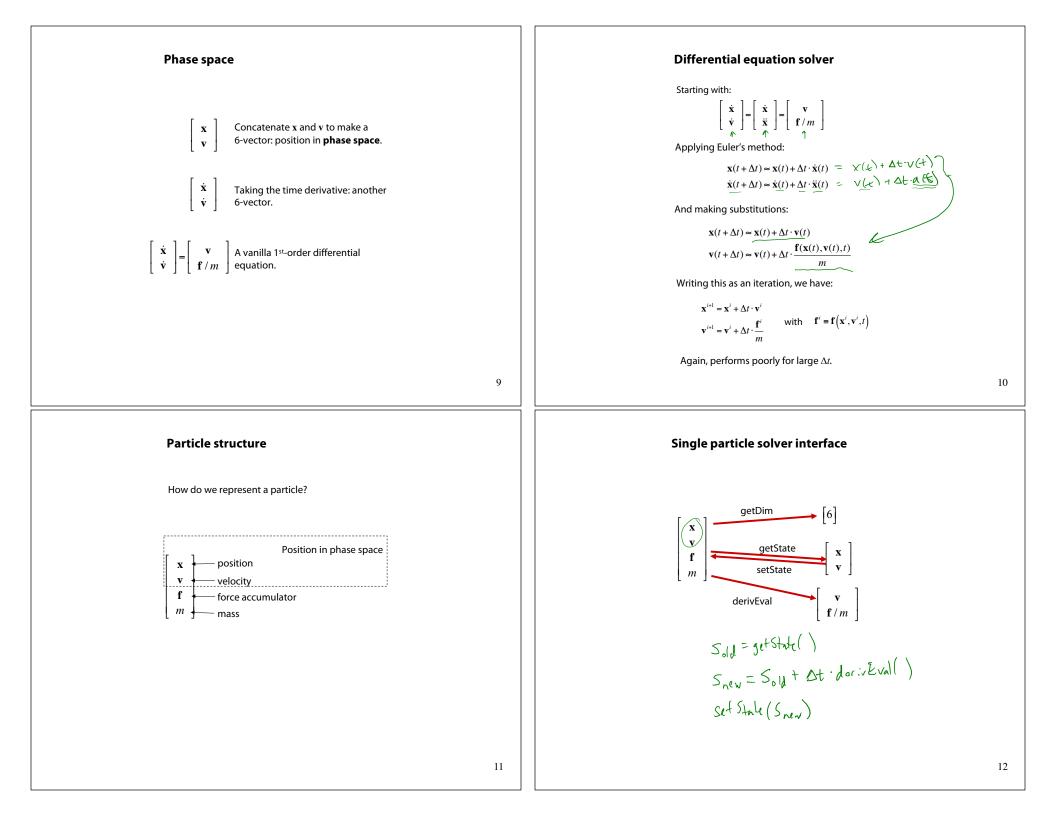
Our solution method, though, worked on first order differential equations.

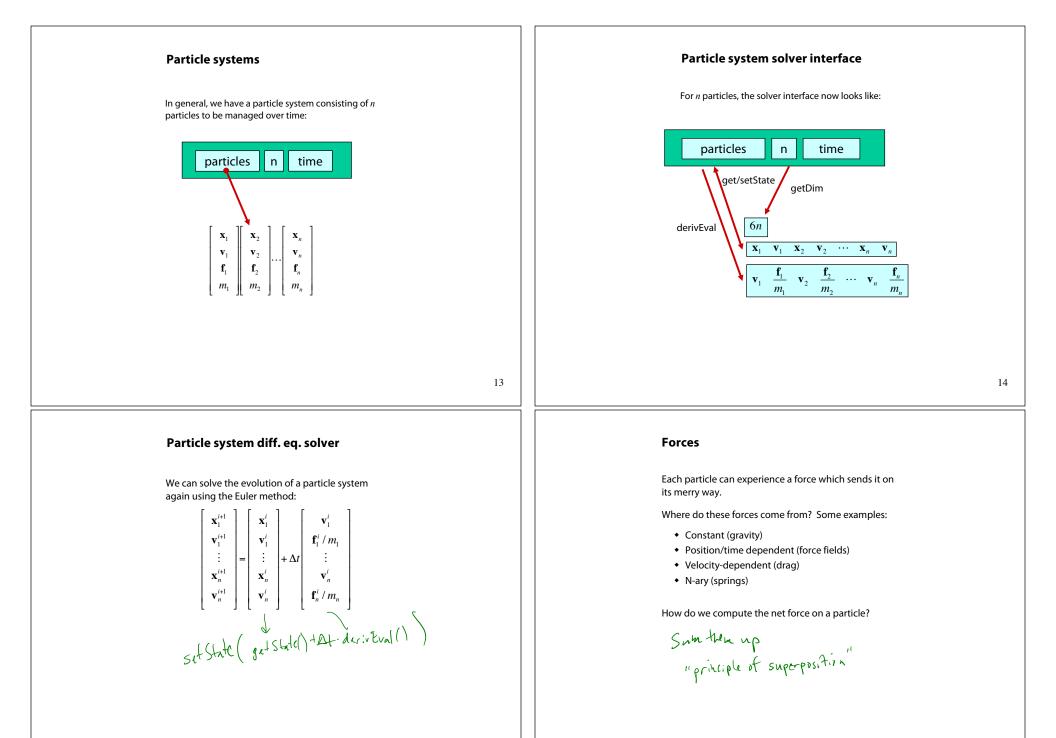
We can rewrite the second order equation as:

$$\begin{bmatrix} \dot{\mathbf{x}} = \mathbf{v} \\ \dot{\mathbf{v}} = \frac{\mathbf{f}(\mathbf{x}, \mathbf{v}, t)}{m} \end{bmatrix} \text{ or } \begin{bmatrix} \dot{\mathbf{x}} \\ \dot{\mathbf{v}} \end{bmatrix} = \begin{bmatrix} \mathbf{v} \\ \frac{\mathbf{f}(\mathbf{x}, \mathbf{v}, t)}{m} \end{bmatrix}$$

where we substitute in v and its derivative to get a pair of **coupled first order equations**.

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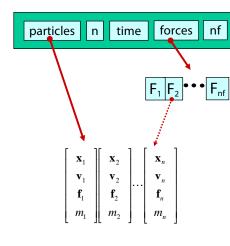




Particle systems with forces

Force objects are black boxes that point to the particles they influence and add in their contributions.

We can now visualize the particle system with force objects:



Damped spring

A spring is a simple examples of an "N-ary" force. Recall the equation for the force due to a 1D spring:

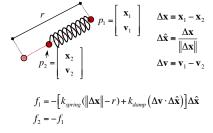
r = rest length

With damping:

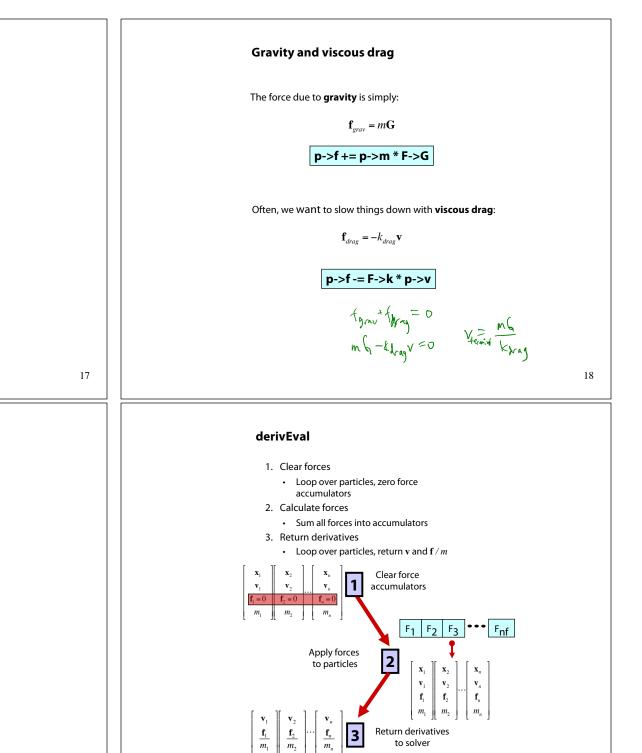
$$f = -[k_{spring}(x - r) + k_{damp}v]$$

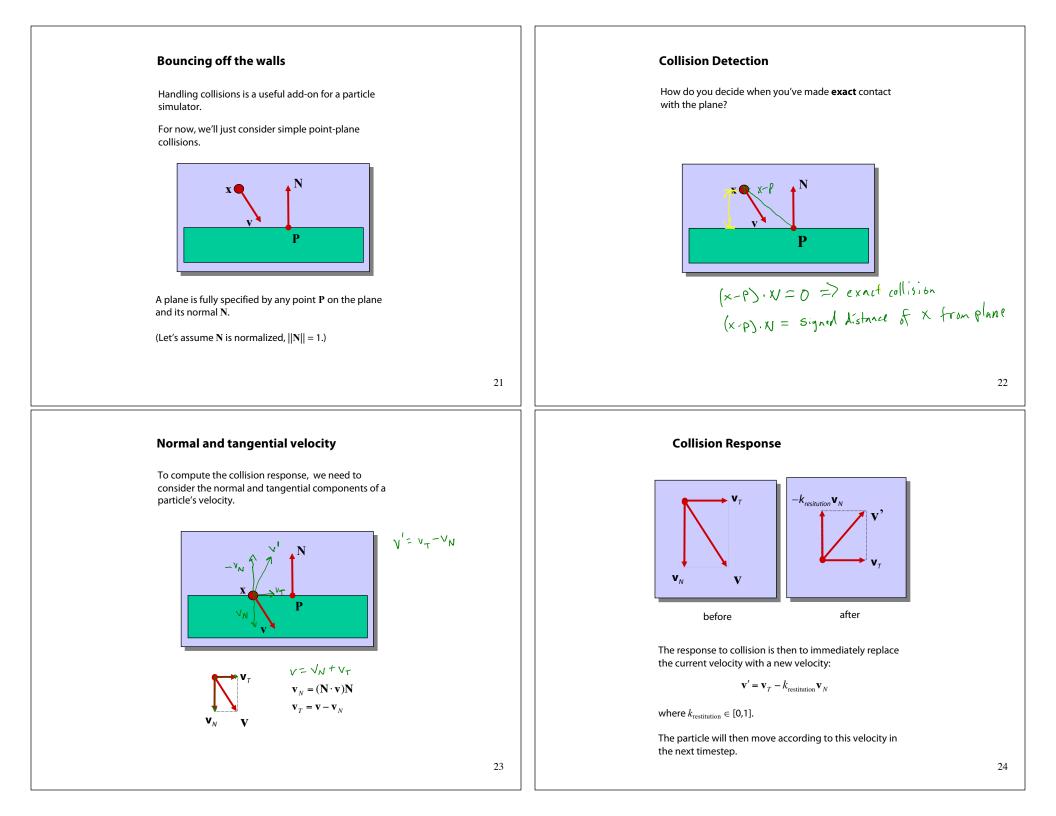
 $f = -k_{spring}(x - r)$

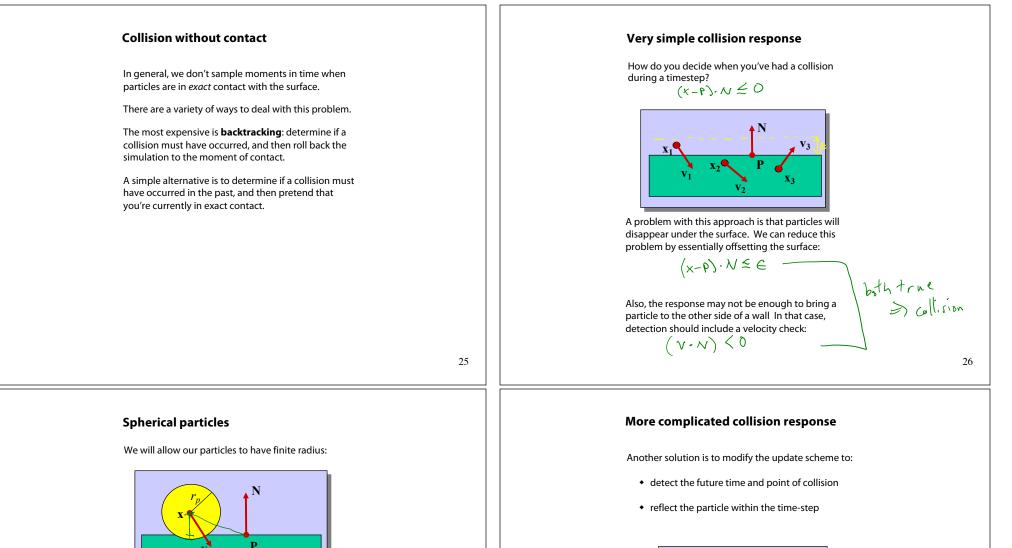
In 2D or 3D, we get:

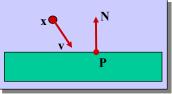


Note: stiff spring systems can be very unstable under Euler integration. Simple solutions include heavy damping (may not look good), tiny time steps (slow), or better integration (Runge-Kutta is straightforward).









The velocity test is unchanged.

 $(x-P)\cdot N \leq r_{P}$

the surface before being reflected: $(x - p) \cdot y \leq f_p + \xi$

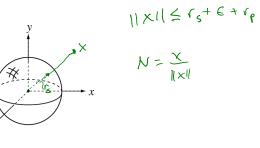
The basic particle-inside-surface test then becomes:

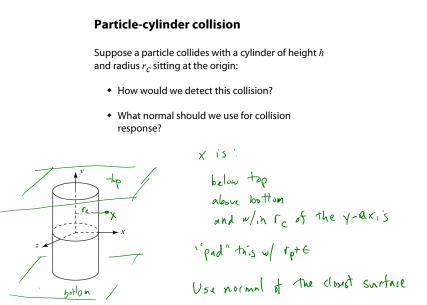
We still pad this test to limit how much particles cross

Particle-sphere collision

Suppose a particle collides with a sphere of radius r_s sitting at the origin:

- How would we detect this collision?
- What normal should we use for collision response?



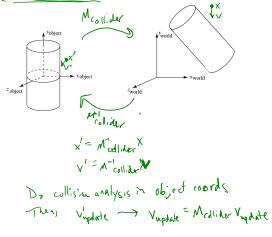


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Collision in model coordinates

As with ray tracing, it convenient to handle collisions in canonical object coordinates.

For the project, we will assume only rigid transformations <u>– rotations+translations</u> – of colliders and particles.



Collision in model coordinates (cont'd)

Note: we assume for the project that scaling will be handled only through the provided dimensions of colliders and particles (height, radius, etc.).

If uniform scale is additionally applied, then it must be extracted from the collider transformation and applied to the epsilons and particle sphere radius when operating in the collider's object coordinates.

If non-uniform scale is applied, then things get trickier, since, e.g., spherical particles may become oriented (nonaxis-aligned) ellipsoids, requiring more complex collision analysis.

Extra credit... 😊

