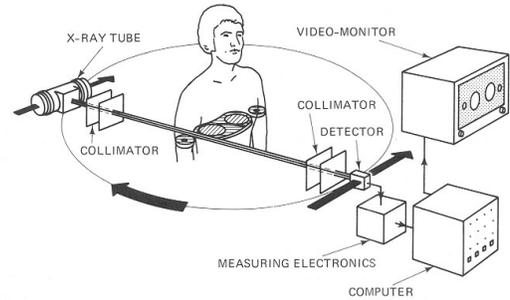


CT Reconstruction

1

Computerized tomography

Computerized tomography (CT) is a method for using x-ray images to reconstruct a spatially varying density function.



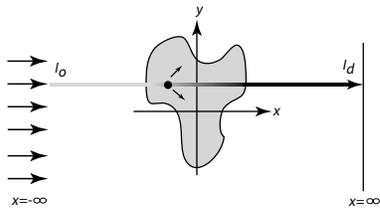
First generation CT scanner

2

Physics of beam attenuation

CT works by collecting x-ray images one slice at a time.

Consider a parallel beam of x-rays passing through an object being imaged orthographically:



An x-ray photon interacts with the material by:

- ♦ absorption
- ♦ scatter

Absorbed photons are simply lost.

We will assume that scattered photons are all re-directed away from the sensor.

3

Physics of beam attenuation

If we consider a single "ray" passing through, we'll find that its intensity drops off as:

$$\Delta I = -\mu I \Delta x$$

where μ is the **linear attenuation coefficient**.

We can re-write this as differentials and permit μ to vary along the ray:

$$dI = -\mu(x)I dx$$

If the material is made of a single substance of varying density, then $\mu(x)$ can be modeled as proportional to that density.

Re-arranging:

$$\frac{dI}{I} = -\mu(x) dx$$

Integrating:

$$\int_{I_0}^{I_d} \frac{dI}{I} = - \int_{-\infty}^{\infty} \mu(x) dx$$

4

Physics of beam attenuation

Performing the integration of the left side:

$$\int_{I_o}^{I_d} \frac{dI}{I} = \ln[I] \Big|_{I_o}^{I_d} = \ln[I_d] - \ln[I_o] = \ln \left[\frac{I_d}{I_o} \right]$$

Equating to the right side:

$$\ln \left[\frac{I_d}{I_o} \right] = - \int_{-\infty}^{\infty} \mu(x) dx$$

Raising to an exponent:

$$\frac{I_d}{I_o} = \exp \left[- \int_{-\infty}^{\infty} \mu(x) dx \right]$$

Solving for detector intensity:

$$I_d = I_o \exp \left[- \int_{-\infty}^{\infty} \mu(x) dx \right]$$

Considering beams that pass through at various y -positions:

$$I_d(y) = I_o \exp \left[- \int_{-\infty}^{\infty} \mu(x, y) dx \right]$$

5

Physics of beam attenuation

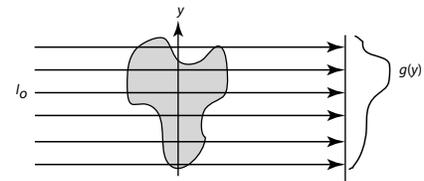
If we back up a little bit, we can remove the negative sign by inverting the argument of the log:

$$\ln \left[\frac{I_o}{I_d} \right] = \int_{-\infty}^{\infty} \mu(x) dx$$

Allowing y to vary:

$$\ln \left[\frac{I_o}{I_d(y)} \right] = \int_{-\infty}^{\infty} \mu(x, y) dx = g(y)$$

Thus, we can take the detector data, and, using this log, we can interpret the result as an integral projection of the attenuation function.

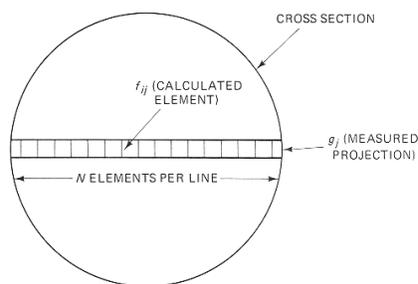


6

ART

Using projections from multiple angles, you can try to solve for the interior distribution.

One approach is essentially to create a large linear system and solve iteratively.



Such a technique is called an Algebraic Reconstruction Technique, or ART.

7

ART

For example:

$$\begin{array}{cccc} 7 & 11 & 9 & 13 \\ & | & | & | \\ & a & b & -12 \\ & | & | & | \\ & c & d & -8 \end{array} \quad \begin{bmatrix} 1 & 0 & 0 & 1 \\ 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 0 & 1 & 1 & 0 \\ 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} = \begin{bmatrix} 7 \\ 11 \\ 9 \\ 13 \\ 12 \\ 8 \end{bmatrix}$$

In practice, ART has proven computationally expensive and sensitive to noise.

Instead, we can use some fancier math to derive an elegant solution...

8

The 1D Fourier transform

Recall (from CSE 557?) that the Fourier transform of a 1D function can be written as:

$$\mathcal{S}_{1D}\{f(x)\} = \int_{-\infty}^{\infty} f(x) \exp[-i2\pi ux] dx = F(u)$$

where u is spatial frequency.

The inverse Fourier transform is simply:

$$\mathcal{S}_{1D}^{-1}\{F(u)\} = \int_{-\infty}^{\infty} F(u) \exp[i2\pi ux] du = f(x)$$

Note that an $f(x)$ implies a unique $F(u)$ and vice versa, so if we know one, we can compute the other:

$$f(x) \xrightarrow{\mathcal{S}} F(u)$$

$$f(x) \xleftarrow{\mathcal{S}^{-1}} F(u)$$

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The 2D Fourier transform

We can generalize this to 2D:

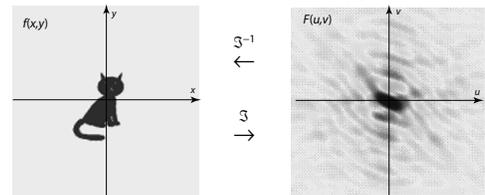
$$\mathcal{S}_{2D}\{f(x, y)\} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) \exp[-i2\pi(ux + vy)] dx dy = F(u, v)$$

where u is spatial frequency in x , and v is the spatial frequency in y .

Likewise, the 2D inverse Fourier transform is:

$$\mathcal{S}_{2D}^{-1}\{F(u, v)\} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(u, v) \exp[i2\pi(ux + vy)] du dv = f(x, y)$$

Again, given one function, we can uniquely compute the other.



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Linear transforms of Fourier domains

We can also write the Fourier transform relation in terms of vector arguments:

$$f(\mathbf{x}) \xrightarrow{\mathcal{S}} F(\mathbf{u})$$

It's easy to show that scaling one domain corresponds to inverse scaling the other:

$$f(a\mathbf{x}) \xrightarrow{\mathcal{S}} \frac{1}{|a|} F\left(\frac{\mathbf{u}}{|a|}\right)$$

In fact, if we replace "a" with a matrix "A", it is not hard to show that:

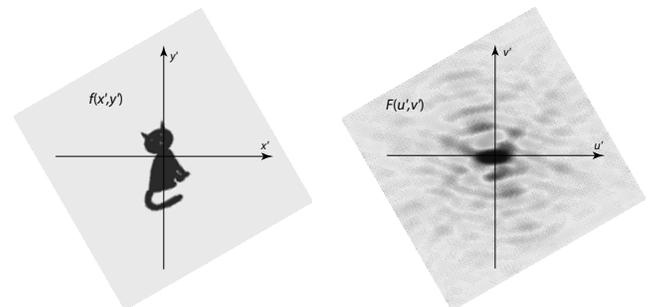
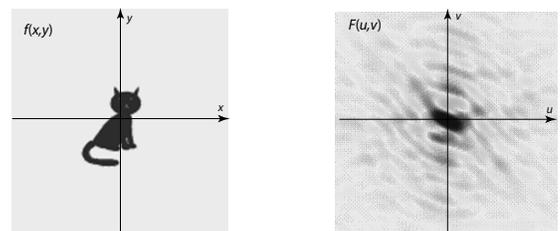
$$f(A\mathbf{x}) \xrightarrow{\mathcal{S}} \left\|A^{-T}\right\| F(A^{-T}\mathbf{u})$$

For rotations, this implies:

$$f(R\mathbf{x}) \xrightarrow{\mathcal{S}} \left\|R^{-T}\right\| F(R^{-T}\mathbf{u}) = ?$$

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Linear transforms of Fourier domains



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Fourier transforms and projections

So, what do Fourier transforms have to do with x-ray projections?

Let's change terminology slightly and say $f(x,y) = \mu(x,y)$. We've already noted that:

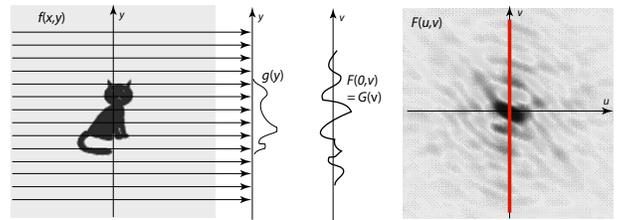
$$F(u,v) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y) \exp[-i2\pi(ux + vy)] dx dy$$

What happens if we evaluate this at $F(0,v)$?

$$\begin{aligned} F(0,v) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y) \exp[-i2\pi(u \cdot 0 + vy)] dx dy \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y) \exp[-i2\pi vy] dx dy \\ &= \int_{-\infty}^{\infty} \left\{ \int_{-\infty}^{\infty} f(x,y) dx \right\} \exp[-i2\pi vy] dy \\ &= \int_{-\infty}^{\infty} g(y) \exp[-i2\pi vy] dy \\ &= \mathfrak{F}_{1D}\{g(y)\} \end{aligned}$$

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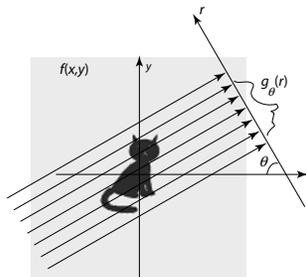
Fourier transforms and projections



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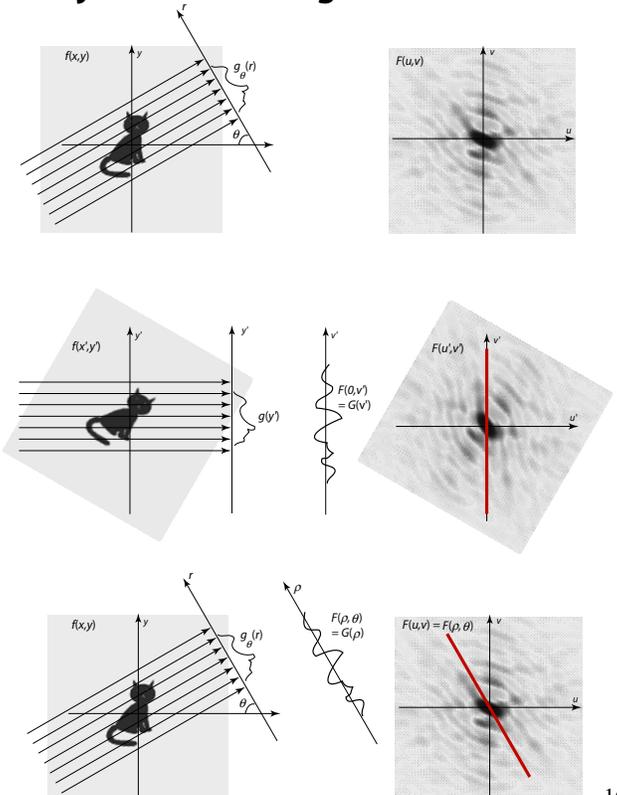
Projection at an angle

What happens if we project the volume at an angle?



15

Projection at an angle



16

Fourier projection slice theorem

In other words, if we express $F(u, v)$ in polar coordinates $F(\rho, \theta)$:

$$F(\rho, \theta) = \mathfrak{F}_{1D}\{g_\theta(r)\} = G_\theta(\rho)$$

This result is called the “Fourier projection slice theorem” or the “central slice theorem.”

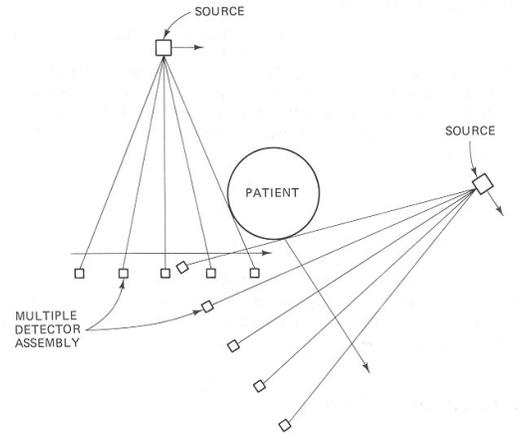
Using this theorem, we can reconstruct an object from its projections by:

1. Populating the Fourier domain with oriented Fourier lines
2. Taking the inverse Fourier transform

In practice, all of these operations can be performed in the spatial domain.

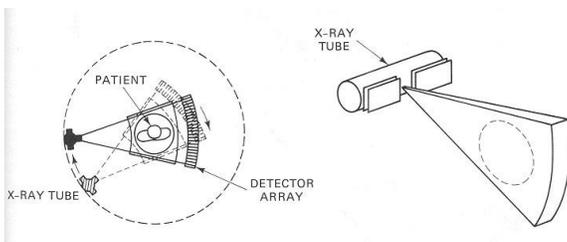
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Second generation scanner



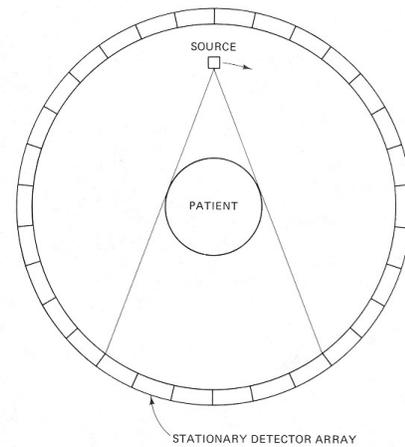
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Third generation scanner



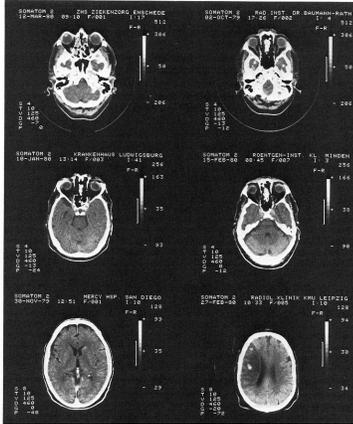
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Fourth generation scanner



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A medical scanner



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