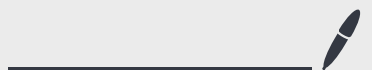


Lecture 19

Dec 3, 2024



Today: Toric Code Correction & Hamiltonian Simulation.

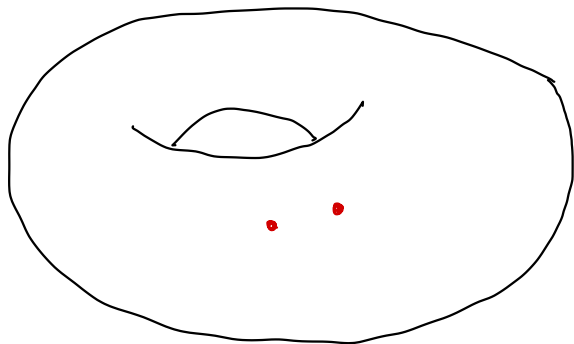
Correction: After measuring the syndrome (i.e. all the X-checks and Z-checks), we want to find the correction that takes us back to the codespace.

Classical repetition code:

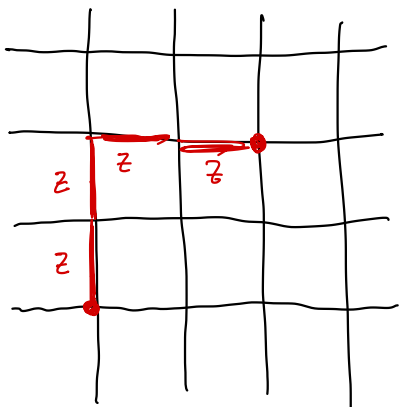
If most bits are 0, correct to 0^n .

If most bits are 1, correct to 1^n .

First toric code example:



Assume syndrome measurements were +1 except at these two vertex checks.



X-checks correct

Z-errors.

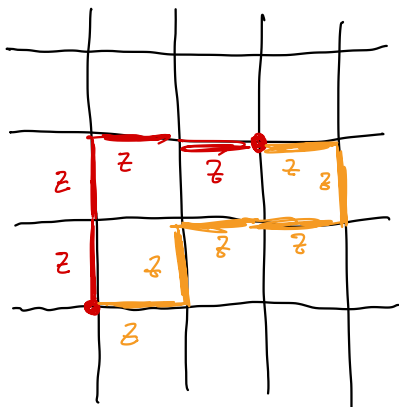
Pauli
 What errors can
 have this syndrome pattern?

Up to a trivial error, it must be a path
 connecting the two end points.

How do we know which path?

We don't. All paths are equivalent up to a
 trivial error.

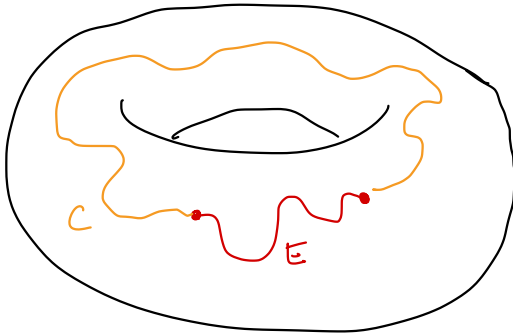
If the true error was in
 red, and the bad error
 in orange, connecting by
 flipping orange,



yields an overall trivial Z -flip as it creates a boundary.

So, short answer is it didn't matter, as long as Error E + Correction C don't generate a non-trivial cycle.

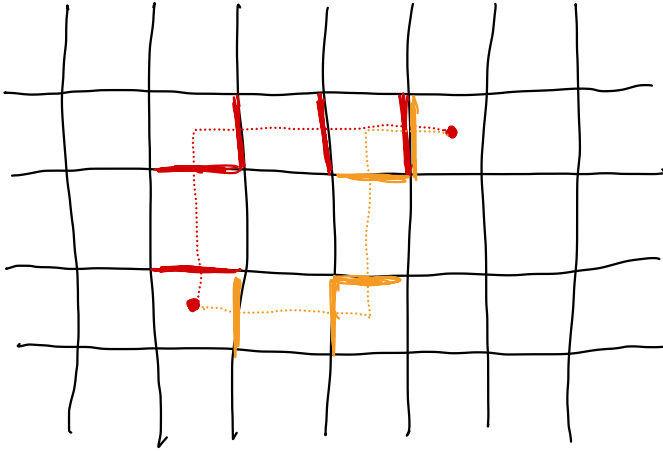
Ex.



If our correction went around the torus, we would have a problem.

Correction assumes errors of size $< O(\sqrt{n})$, so we should only find corrections that are also short.

Note: computation of the correction can be done from the syndrome by a classical computer.



Hamiltonian Simulation

One of the strongest applications of q. computers is simulating q. mechanics which we believe would take exponential time on a classical computer.

Given a Hamiltonian $H = H(t)$, the evolution of a state is

defined by
$$\frac{d}{dt} |\psi(t)\rangle = H(t) |\psi(t)\rangle.$$

When H is time-invariant,

$$|\psi(t)\rangle = e^{-iHt} |\psi(0)\rangle.$$

Matrix exponentiation: if $H = \sum_a \lambda_a |a\rangle\langle a| \leftarrow$ Spectral decomposition

then $e^{-iHt} = \sum_a e^{-i\lambda_a t} |a\rangle\langle a| \leftarrow$ unitary

Furthermore, $e^{-i(UHU^\dagger)t} = U e^{-iHt} U^\dagger$

PF. $UHU^\dagger = \sum_a \lambda_a \underbrace{|u|a\rangle}_{\substack{\uparrow \\ \text{new eigenvectors}}} \langle a|U^\dagger$

$$\begin{aligned} e^{-i(UHU^\dagger)t} &= \sum_a e^{-i\lambda_a t} |u|a\rangle \langle a|U^\dagger \\ &= U \left(\sum_a e^{-i\lambda_a t} |a\rangle \langle a| \right) U^\dagger. \quad \square \end{aligned}$$

If we can diagonalize a Hamiltonian H s.t. $H = U H_{\text{diag}} U^\dagger$

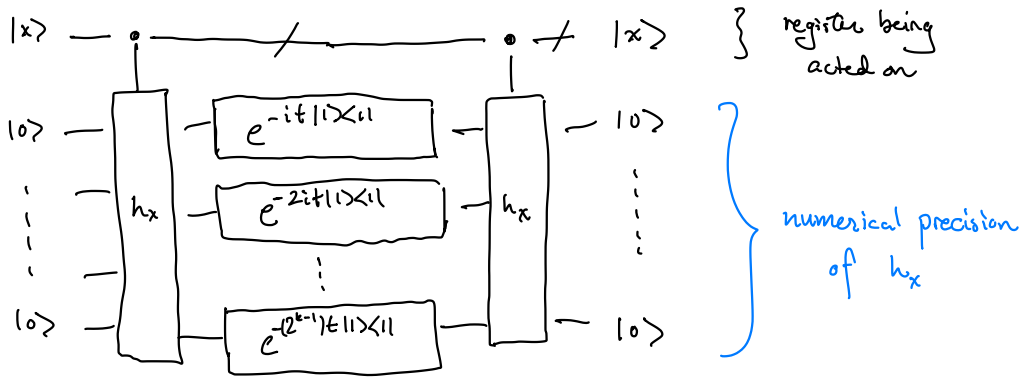
then evolving w.r.t. H involves 2 applications of U +

evolving by H_{diag} .

$$H_{\text{diag}} = \sum_x h_x |x\rangle\langle x| = \begin{pmatrix} h_0 & & \\ & \dots & \\ & & h_{2^L-1} \end{pmatrix}$$

↑
real.

Circuit for $e^{-itH_{diag}}$



↑ write h_x as a number

$$\text{If } h_x = h_x^{(0)} + 2h_x^{(1)} + 4h_x^{(2)} + \dots + (2^{k-1})h_x^{(k-1)}$$

binary decomposition, then we map

$$\begin{aligned} |x\rangle &\mapsto e^{-ih_x^{(0)}t} \cdot e^{-2ih_x^{(1)}t} \cdot \dots \cdot e^{-2^{k-1}ih_x^{(k-1)}t} |x\rangle \\ &= e^{-ih_x t} |x\rangle \end{aligned}$$

When H_{diag} is local, acting on a few qubits, this is implementable.

Furthermore, local Hamiltonian terms h_j can be diagonalized

easily. Therefore, we can implement the evolution of local Hamiltonian terms h_j easily.

What about the evolution of $H = \sum h_j$?

Even though each h_j is diagonalizable, the net sum may be hard to diagonalize and would end up a matrix on n -qubits

Solution:

$$e^{-i(h_1+h_2)t} = \lim_{m \rightarrow \infty} \left(e^{-\frac{ih_1 t}{m}} e^{-\frac{ih_2 t}{m}} \right)^m$$

Lie product formula.

$$\left\| \left(e^{-\frac{ih_1 t}{m}} e^{-\frac{ih_2 t}{m}} \right)^m - e^{-i(h_1+h_2)t} \right\| \leq \epsilon$$

$$\text{for } m = O\left(\frac{(vt)^2}{\epsilon}\right) \text{ where } v = \max\{\|h_1\|, \|h_2\|\}.$$

Unappealing to evolve t^2 time for t -evolution, but there are speedups to get it to $t^{1+\delta}$ for any $\delta > 0$.

general evolution of $H = \sum_{j=1}^J h_j$. Let $m = O\left(\frac{(vt)^2}{\epsilon}\right)$

where $v = \max\{\|h_1\|, \dots, \|h_J\|\}$ (typically 1 wlog).

① Compute circuit for $e^{-ih_j t'}$ for $t' = t/m$.

using diagonalization + diagonal evolution

② apply $\left(e^{-ih_1 t'} \quad e^{-ih_2 t'} \quad \dots \quad e^{-ih_J t'} \right)^m$.

Runtime = $\frac{Jm^2}{\epsilon}$ and can be improved to $\frac{Jm^{1.8}}{\epsilon}$.

Other simulable Hamiltonians include sparse Hamiltonians

where H is mostly 0 except $\text{poly}(n)$ entries per row.

\exists efficient alg for simulating evolution in this case.