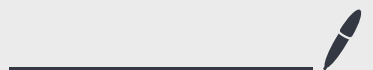


Lecture 12

Nov 5, 2024



Continuation of prev lecture on Shor's:

Only remains to show with prob.  $\geq \frac{1}{32}$ ,

$$-\frac{r}{2} \leq yr \bmod Q \leq \frac{r}{2}. \quad (*)$$

With this, we conclude, we generate  $\text{ord}(x)$  with probability,

$$\Omega\left(\frac{1}{\log N}\right).$$

Pf of (\*):

Recall amplitude on  $|\gamma\rangle$  is  $\frac{1}{\sqrt{QJ}} \omega^{\gamma r} \sum_{j=0}^{J-1} \omega^{\gamma r j}$

with  $J = \lfloor \frac{Q}{r} \rfloor$

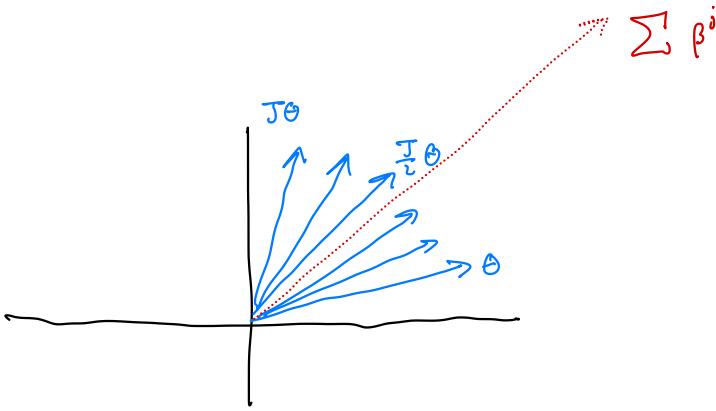
Focus on the  $\sum_{j=0}^{J-1} \omega^{\gamma r j}$  term as this is where the constructive/  
destructive interference occurs. Let  $\beta = \omega^{\gamma r} = e^{(2\pi i \cdot \frac{\gamma r}{Q})}$

If  $-\frac{r}{2} \leq \gamma r \pmod{Q} \leq \frac{r}{2}$ , then  $\beta = e^{i\theta}$  for  $\theta$  an  
angle s.t.  $|\theta| \leq \frac{2\pi r}{2Q} = \pi \left(\frac{r}{Q}\right)$ .

Then  $\omega^{\gamma r j} = \beta^j$  corresponds to angle  $j\theta$  with

$$|j\theta| \leq |J\theta| \leq \pi.$$

So, the terms of  $\sum_{j=0}^{J-1} \omega^{\gamma r j} = \sum_{j=0}^{J-1} \beta^j$  span only angle  $\pi$ .



Simple calculation:  $\frac{1}{2}$  the terms make angle  $\leq \frac{\pi}{4}$  to resultant vector. Since overall span  $\leq \pi$ , no vector contributes negatively to resultant vector.

$$\cos \frac{\pi}{4} = \frac{1}{\sqrt{2}}.$$

So, length of resultant vector of  $\frac{1}{\sqrt{QJ}} \omega^{\gamma \ell} \sum_{j=0}^{J-1} \omega^{\gamma r_j}$

$$\geq \frac{J}{2} \cdot \frac{1}{\sqrt{2}} \cdot \frac{1}{\sqrt{QJ}} = \frac{1}{2^{3/2}} \cdot \frac{1}{\sqrt{\frac{Q}{J}}} \geq \frac{1}{4\sqrt{r}}.$$

$\uparrow$  half terms  
 $\uparrow$  contribution per term

Conclusion: If  $-\frac{r}{2} \leq \gamma r \pmod{Q} \leq \frac{r}{2}$ , then

the resulting vector  $|\gamma\rangle$  has magnitude  $\geq \frac{1}{4\sqrt{r}}$ .

We next show many such vectors  $\gamma$  exist.

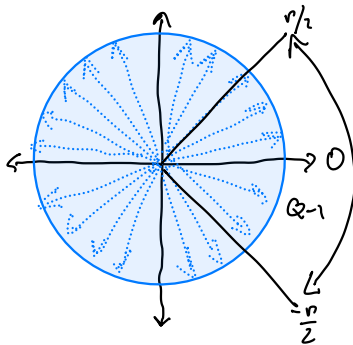
(since  $r \in \mathbb{Z}_{\mathbb{Q}}^*$ )

If  $\gcd(r, Q) = 1$ , then  $\exists r^{-1}$  s.t.  $r \cdot r^{-1} \equiv 1 \pmod{Q}$ .

Therefore the map  $\gamma \mapsto \gamma r$  is a permutation of  $\{0, \dots, Q-1\}$ .

So, at least  $r$  vectors  $\gamma$  exist. s.t.  $-\frac{r}{2} \leq \gamma r \pmod{Q} \leq \frac{r}{2}$ .

Cartoon:



all possible values occur as  $\gamma r$ .

So  $r$  vectors in region.

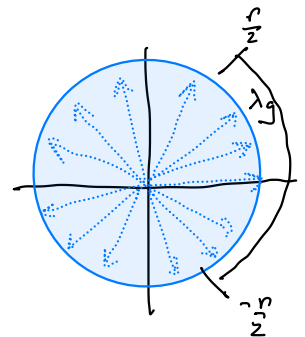
When  $g := \gcd(r, Q) > 1$ , note

$g \leq r/2$ . Then  $\gamma r$  are

uniformly distributed over

$$0, g, 2g, \dots, \left(\frac{Q}{g} - 1\right)g$$

with  $\gamma r = \lambda g$  for  $g$  values of  $\gamma$ .



$$\text{If } |\lambda g| < \frac{r}{2} \Rightarrow |\lambda| \leq \frac{r}{2g}$$

Then, for at least  $2 \left\lfloor \frac{r}{2g} \right\rfloor \cdot g \geq \frac{r}{2}$  vectors  $\gamma_i$

$$-\frac{r}{2} \leq \gamma r \pmod{Q} \leq \frac{r}{2}.$$

So, total probability mass on  $\gamma$  s.t.  $-\frac{r}{2} \leq \gamma r \pmod{Q} \leq \frac{r}{2}$ .

$$\text{is } \geq \frac{r}{2} \cdot \left( \frac{1}{4\sqrt{r}} \right)^2 = \frac{1}{32}.$$


Therefore, we sample a  $\gamma$  according to the algorithm for order finding, we correctly calculate  $\text{ord}(x)$  with probability  $\Omega\left(\frac{1}{\log N}\right)$ .

So our overall algorithm for factoring is efficient in that it runs in time  $\text{polylog}(N)$ .

Next time: efficient classical algorithms for simulating quantum computation.

# Today: Efficient classical algorithms for simulating quantum computations

Problem: given an input  $\langle C \rangle$  the description of a q. circuit with  $n$  qubits,  $m$  gates, and no measurements,

what is the probability that   $\} |0^n\rangle$

the measurement output is 1?

$w =$  mat. mult. coefficient.

Computing  $p = \|\langle 1 | \otimes \mathbb{1} | C | 0^n \rangle\|^2$  to accuracy  $\epsilon$  can be solved in classical time  $O(2^{wn} \log(\frac{m}{\epsilon}))$  and space  $O(2^n \log(\frac{m}{\epsilon}))$ .

Pf. Let  $C = g_m g_{m-1} \dots g_1$ .

For gate  $g_t$ , let  $\tilde{g}_t$  be the pruning of  $g_t$  to  $\ell$  bits.

Then  $\|\tilde{g}_t - g_t\|_F = \sqrt{\sum_{ij} (\tilde{g}_t - g_t)_{ij}^2} = 4 \cdot 2^{-\ell}$  for  $4 \times 4$  matrices.

Since  $\|\cdot\| \leq \|\cdot\|_F$ ,

$$\|\tilde{g}_\ell - g_\ell\| \leq 4 \cdot 2^{-\ell} \quad \text{so} \quad \|\tilde{g}_\ell \otimes \mathbb{1} - g_\ell \otimes \mathbb{1}\| \leq 4 \cdot 2^{-\ell}.$$

Then for  $\tilde{C} = \tilde{g}_m \tilde{g}_{m-1} \dots \tilde{g}_1$ ,

$$\|\tilde{C}|0^n\rangle - C|0^n\rangle\| \leq 4m \cdot 2^{-\ell}$$

$$\text{So } |\tilde{p} - p| \leq 8m \cdot 2^{-\ell} \leq \epsilon$$

Choose  $\ell$  s.t.  $8m \cdot 2^{-\ell} \leq \epsilon \Rightarrow \ell \geq \Omega\left(\log \frac{m}{\epsilon}\right)$ .

Compute  $\tilde{p}$  using mat. multiplication.  $|p - \tilde{p}| \leq \epsilon$ .

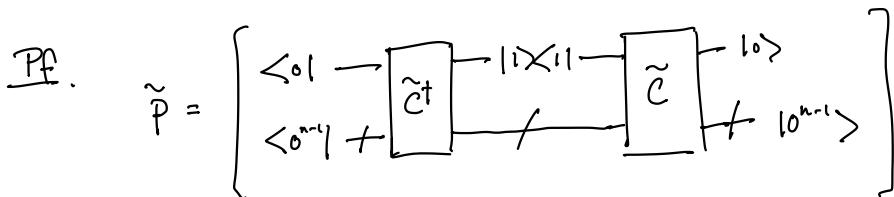
Additionally, we can multiply and prune as we compute.

gives a runtime of  $O\left(2^{wn} \log\left(\frac{m}{\epsilon}\right)\right)$  and space  $O\left(2^n \log\left(\frac{m}{\epsilon}\right)\right)$ .

$\uparrow$                      $\uparrow$   
 mat. mult.       size of  
                          integers

In actuality, no one uses such fast mult. algorithms since the coefficients are huge. So runtime is more like  $O\left(2^{2.77n} \log^2\left(\frac{m}{\epsilon}\right)\right)$ .

Claim We can reduce the space complexity to  $\text{poly}\left(n, \log\left(\frac{1}{\epsilon}\right)\right)$ .





$$\tilde{p} = \langle 0^n | \tilde{g}_1^\dagger \tilde{g}_2^\dagger \dots \tilde{g}_m^\dagger (|1\rangle\langle 1| \otimes \mathbb{1}) \tilde{g}_m \dots \tilde{g}_1 | 0^n \rangle$$

(one big matrix multiplication)

Add identity terms  $\mathbb{1} = \sum_{\gamma \in \{0,1\}^n} |\gamma\rangle\langle\gamma|$ .

$$\begin{aligned} \tilde{p} &= \langle 0^n | \tilde{g}_1 \left( \sum_{\gamma_{2m+1}} |\gamma_{2m+1}\rangle\langle\gamma_{2m+1}| \right) \tilde{g}_2 \left( \sum_{\gamma_{2m}} |\gamma_{2m}\rangle\langle\gamma_{2m}| \right) \dots \left( \sum_{\gamma_1} |\gamma_1\rangle\langle\gamma_1| \right) \tilde{g}_1 | 0^n \rangle \\ &= \sum_{\gamma_1, \dots, \gamma_{2m+1}} \langle 0^n | \tilde{g}_1 | \gamma_{2m+1} \rangle \dots \langle \gamma_1 | \tilde{g}_1 | 0^n \rangle. \end{aligned}$$

Alg: Iterate over  $\gamma_1, \dots, \gamma_{2m+1} \in \{0,1\}^n$  computing each multiplication in the sum. Requires

$$O\left(2^{2nm} \log^2\left(\frac{m}{\epsilon}\right)\right) \text{ time but only } O(nm + \log\left(\frac{m}{\epsilon}\right)) \text{ space.}$$

To estimate  $p$  to  $1/\epsilon$ , requires only  $O(nm)$  space.

Proves  $BQP \subseteq PSPACE$ . (Called the Feynman path integral)

i.e. every  $q$ . computation can be simulated with polynomial space but (perhaps) exponential time.

Next: A situation when we can vastly improve the time complexity.

The issue is that keeping track of the state  $g_1 \dots g_n |0^n\rangle$  is inefficient and may take  $2^n$  complex numbers to record.

One solution was to keep "no numbers" using path integral.

Another is succinct descriptions of  $g_i$  states.

First, Pauli matrices:

$$\mathbb{1}, X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = iXZ, Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

$X, Y, Z$  all anticommute, have trace 0, square to  $\mathbb{1}$ .

$$\mathcal{P}_1 = \left\{ \pm \mathbb{1}, \pm i\mathbb{1}, \pm X, \pm iX, \pm Y, \pm iY, \pm Z, \pm iZ \right\}.$$

a group under matrix multiplication.

$$\mathcal{P}_n = \left\{ P_1 \otimes P_2 \otimes \dots \otimes P_n \mid P_1, \dots, P_n \in \mathcal{P}_1 \right\}. \quad \text{Also a group.}$$

Pauli matrices can be described with  $2(n+1)$  bits.

Use notation:  $X_j$  to denote  $\mathbb{1} \otimes \dots \otimes \mathbb{1} \otimes X \otimes \mathbb{1} \otimes \dots \otimes \mathbb{1}$   
 $\uparrow$   
 $j^{\text{th}}$  location.

$$\begin{aligned} \text{So } (X_1 Z_2)(X_2 Y_3) &= (X \otimes Z \otimes \mathbb{1})(\mathbb{1} \otimes X \otimes Y) \\ &= X \otimes Z X \otimes Y \\ &= X \otimes iY \otimes Y = i(X \otimes Y \otimes Y) \\ &= i X_1 Y_2 Y_3. \end{aligned}$$

An observation:  $|0^n\rangle$  is the unique solution to  $Z_j |\psi\rangle = |\psi\rangle$ ,  
 for all  $j = 1, \dots, n$ .

Pf.  $Z_j |\psi\rangle = |\psi\rangle$  only if  $|\psi\rangle = |0\rangle \otimes |\psi'\rangle$ .

Rest follows similarly.  $\square$

Another observation:  $|+\rangle^{\otimes n}$  is the unique solution to  $X_j |\psi\rangle = |\psi\rangle$ ,  
 for all  $j = 1, \dots, n$ .

Lemma Assume  $|\psi\rangle$  is the unique solution to  $P_j |\psi\rangle = |\psi\rangle$  for  
 Pauli matrices  $P_1, \dots, P_n$ . Let  $U$  be any unitary.

Define  $Q_j = UP_jU^\dagger$ . Then  $U|\psi\rangle$  is the unique solution to  $Q_j|\tau\rangle = |\tau\rangle$  for all  $j=1, \dots, n$ .

PP. To see it is a solution, notice

$$\begin{aligned} Q_j U|\psi\rangle &= UP_jU^\dagger U|\psi\rangle \\ &= UP_j|\psi\rangle \\ &= U|\psi\rangle. \end{aligned}$$

For uniqueness, assume  $\exists$  a solution  $|\tau\rangle$ . Then,

$$|\tau\rangle = Q_j|\tau\rangle \implies U^\dagger|\tau\rangle = P_jU^\dagger|\tau\rangle \quad \forall j=1, \dots, n.$$

So,  $U^\dagger|\tau\rangle = |\psi\rangle$  by uniqueness. So  $|\tau\rangle = U|\psi\rangle$ .  $\square$

If  $|\psi\rangle$  is the unique state s.t.  $P_j|\psi\rangle = |\psi\rangle$  for all  $j=1, \dots, n$ , we say  $P_1, \dots, P_n$  stabilize  $|\psi\rangle$ .

Issue is that for arbitrary  $U$ ,  $UP_jU^\dagger$  may not be a Pauli matrix.

But for some  $U$  it will be. The set of  $U$  for which

$UPU^\dagger$  is also a Pauli matrix  $\forall P$  is called the normalizers

group of  $\mathcal{P}_n$ . The normalizer group of  $\mathcal{P}_n$  is called the Clifford group,  $C_n$ .

$$C_n = \left\{ U \mid U P U^\dagger \in \mathcal{P}_n \ \forall P \in \mathcal{P}_n \right\}.$$

It's a more complicated pf than we have time for this class, but every matrix  $\in C_n$  can be generated from

$CNOT \otimes \mathbb{1}_{n-2}$ ,  $S \otimes \mathbb{1}_{n-1}$ ,  $H \otimes \mathbb{1}_{n-1}$ . and their  $\pm$ ,  $\pm i$  variants.

Here,  $S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$ .

Many other unitaries such as  $X, Y, Z, \text{SWAP}, CZ$  are all part of the Clifford group.

Consider  $H_i$  in the Clifford group  $C_n$ .

Suppose  $P_1 \dots P_n$  stabilize  $|\psi\rangle$ .

Then, we can efficiently calculate stabilizers for  $H_i(|\psi\rangle)$ .

If  $P_j = \mathbb{1} \otimes \text{---}$ , then  $H_1 P_j H_1^\dagger = \mathbb{1} \otimes \text{---} = P_j$ .

If  $P_j = X \otimes \text{---}$ , then  $H_1 P_j H_1^\dagger = Z \otimes \text{---}$

If  $P_j = Z \otimes \text{---}$ , then  $H_1 P_j H_1^\dagger = X \otimes \text{---}$

If  $P_j = Y \otimes \text{---}$ , then  $H_1 P_j H_1^\dagger = Y \otimes \text{---}$

Similar rules can be generated for CNOT and S updates.

### ✓ Gottesman-Knill

Thm given a circuit  $g_m \dots g_1$  with each  $g_k \in \{CNOT, S, H\}$ , we can efficiently compute a collection of stabilizers for  $g_m \dots g_1 |0^n\rangle$ .

Pr. Starting with  $P_j = Z_j$  which stabilize  $|0^n\rangle$ , we update stabilizers gate by gate. Each update takes  $O(n^2)$  time as there are  $n$  stabilizers each of  $O(n)$  bits. Total time is  $O(mn^2)$ , space  $O(n^2)$ .  $\square$

What about measurements?

Wlog, we only need to consider measuring the first qubit in standard basis.

Notice if  $P|\psi\rangle = P'|\psi\rangle = |\psi\rangle$  for Paulis  $P, P'$  then  
 $PP'|\psi\rangle = P|\psi\rangle = |\psi\rangle$  so  $PP'$  stabilizes  $|\psi\rangle$  as well.

So if  $P_1, \dots, P_n$  stabilize  $|\psi\rangle$  then

$\langle P_1, \dots, P_n \rangle$  stabilizes  $|\psi\rangle$  where this is the stabilizer subgroup  $\subseteq \mathcal{P}_n$ .

Let  $S_\psi = \{ P \in \mathcal{P}_n \mid P|\psi\rangle = |\psi\rangle \}$ .

Measuring  $|\psi\rangle$ :

- ① If  $Z_i \in S_\psi$ , then measurement outcome is 0 and state doesn't change. Deterministic measurement.
- ② If  $-Z_i \in S_\psi$ , then measurement outcome is 1 and state doesn't change. Deterministic measurement.
- ③ If  $Z_i \notin S_\psi$ , things get more complicated.

$Z_i$  must not commute with all of  $S_\psi$ .

Find a basis for  $S_\psi$  s.t.  $S_\psi = \langle b_1, \dots, b_n \rangle$ ,  
and  $b_1 Z_i = -Z_i b_1$  but  $b_j Z_i = Z_i b_j$  for  $j > 1$ .

Flip a coin. Replace  $b_1$  with  $Z_1$  or  $-Z_1$  depending on the coin flip.

### Pf of correctness

Since  $b_1$  and  $Z_1$  anticommute, square to  $\mathbb{1}$ , by part 2 problem, there exists a change of basis s.t.

$$U b_1 U^\dagger = X_1 \quad \text{and} \quad U Z_1 U^\dagger = Z_1, \quad \text{and} \quad U b_j U^\dagger = \mathbb{1} \otimes b_j'.$$

Since  $b_1 \in S_\psi$ ,  $U|\psi\rangle = |+\rangle \otimes \text{---}$

So measuring,  $Z_1$  is a coin-flip resulting in  $|0\rangle$  or  $|1\rangle$ .

Doesn't change remainder of state, so new state is

stabilized by  $Z_1, b_2, \dots, b_n$  or  $-Z_1, b_2, \dots, b_n$

depending on outcome.  $\square$

Finding a basis  $\langle b_1, \dots, b_n \rangle$  for  $S_\psi$  s.t. only  $b_1$

anticommutes:

① Renumber bases s.t.  $b_1$  anticommutes.

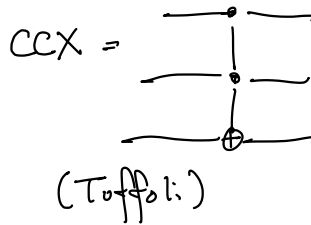
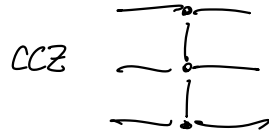


(2) If  $b_k$  anticommutes, replace  $b_k$  with  $b_1 b_k$ .

Next, computation with a few non-Clifford gates.

non-Clifford gate examples:

$$T = \sqrt{S} = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/8} \end{pmatrix}$$



Theorem (Solovay-Kitaev) Any 2-qubit unitary can be  $\epsilon$ -approximated using  $O(\text{polylog}(1/\epsilon))$   $H, T, \text{CNOT}$  gates.

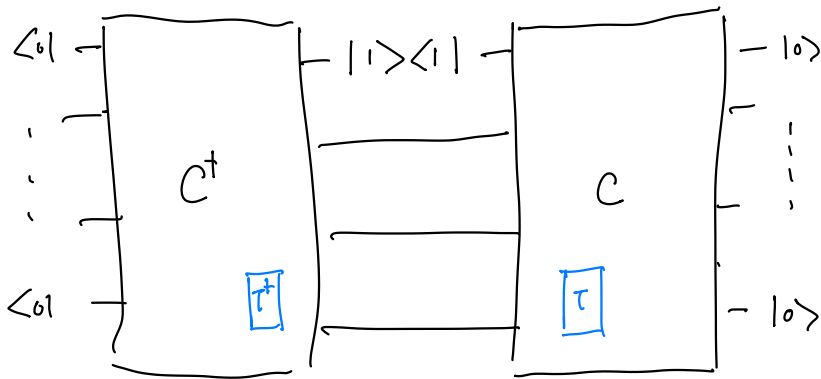
Solovay-Kitaev + Clifford simulation suggests that the number of  $T$  gates in a  $H, T, \text{CNOT}$  circuit should be a measure of the circuit's complexity.

Thm  $\exists$  a constant  $\alpha > 0$ , s.t. computing the output probability of a quantum circuit consisting of  $m$ - Clifford gates,  $t$  T-gates on  $n$  qubits can be classically computed in time  $O(2^{\alpha t} \cdot \text{poly}(n, m))$ .

Best:  $\alpha < 0.9$  (Qassim-Pashyan-Gorriet)

Today  $2^\alpha = 3$ ,  $\alpha < 1.6$ .

Model of such a computation:



$\uparrow$  one big matrix multiplication:

Replacement :

$$T^\dagger \otimes T = a \mathbb{1} \otimes \mathbb{1} + b S^\dagger \otimes S + c Z^\dagger \otimes Z.$$

$$\begin{pmatrix} 1 & & \\ e^{i\pi/4} & & \\ & e^{-i\pi/4} & \\ & & 1 \end{pmatrix} = \begin{pmatrix} a & & \\ & a & \\ & & a \end{pmatrix} + \begin{pmatrix} b & & \\ bi & & \\ & -bi & \\ & & b \end{pmatrix} + \begin{pmatrix} c & & \\ & -c & \\ & & -c \\ & & & c \end{pmatrix}$$

Solve  $a + b + c = 1$

$a = \frac{1}{2}$

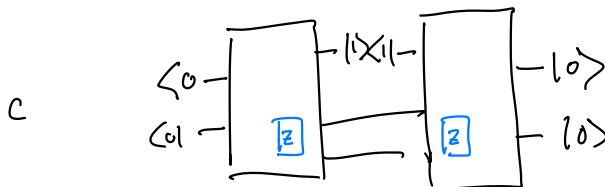
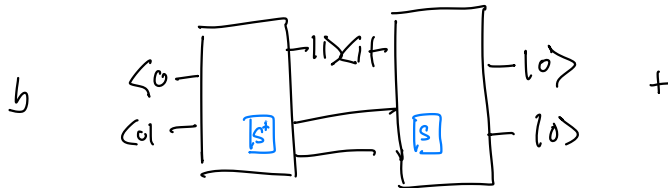
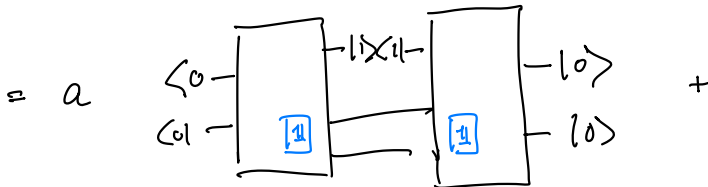
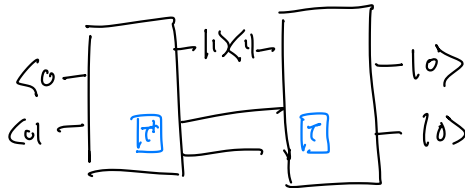
$a + bi - c = e^{i\pi/4}$

$b = \frac{1}{\sqrt{2}}$

$a - bi - c = e^{-i\pi/4}$

$c = \frac{1}{2} - \frac{1}{\sqrt{2}}$

By linearity,



Apply this replacement recursively for every pair of T gates.

Yields  $3^t$  calculations each of which was a only Clifford computation. So previous, subroutine gives an efficient  $\text{poly}(n, m)$  algorithm.