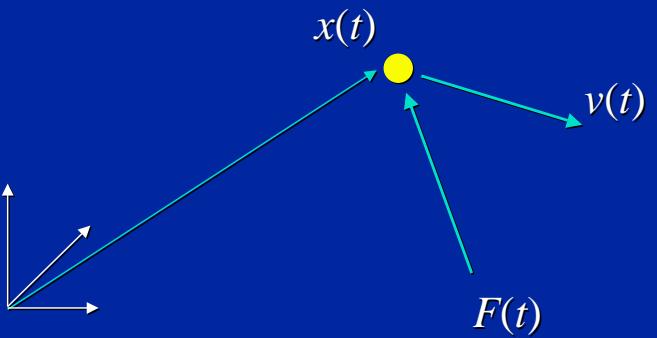


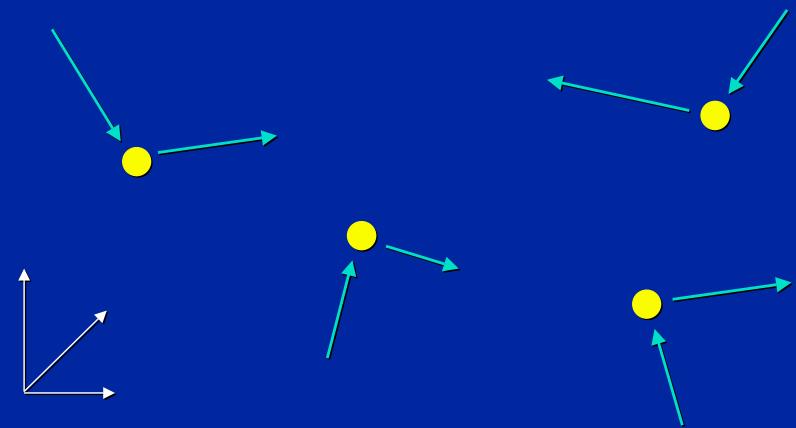
## Rigid Body Simulation



## Particle State

$$\mathbf{X} = \begin{pmatrix} x(t) \\ y(t) \end{pmatrix}$$
$$\mathbf{X} = \underbrace{\begin{array}{|c|c|c|c|c|c|}\hline & x(t) & & v(t) & & \\ \hline & \boxed{\phantom{0}} & \boxed{\phantom{0}} & \boxed{\phantom{0}} & \boxed{\phantom{0}} & \boxed{\phantom{0}} \\ \hline \end{array}}_{\text{State Vector}}$$

## Particle Dynamics



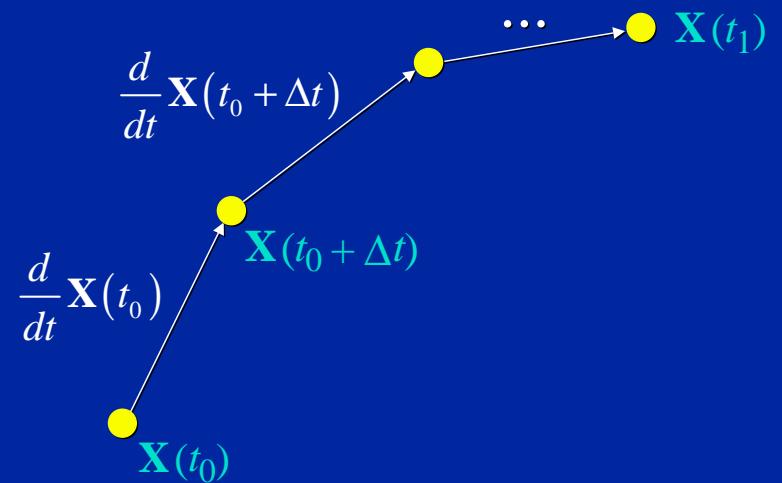
## Multiple Particles

## State Derivative

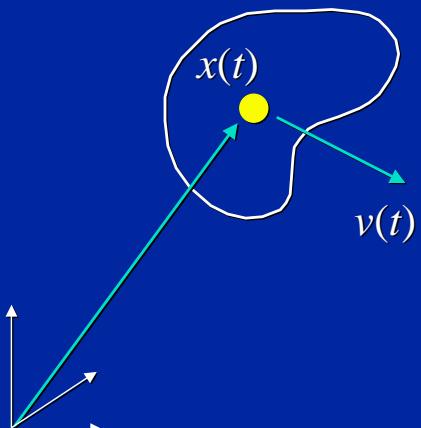
$$\frac{d}{dt} \mathbf{X} = \frac{d}{dt} \begin{pmatrix} x_1(t) \\ v_1(t) \\ \vdots \\ x_n(t) \\ v_n(t) \end{pmatrix} = \begin{pmatrix} v_1(t) \\ F_1(t)/m_1 \\ \vdots \\ v_n(t) \\ F_n(t)/m_n \end{pmatrix}$$

$$\frac{d}{dt} \mathbf{X} = \boxed{\quad} \boxed{\quad} \dots \text{ 6n elements } \dots \boxed{\quad} \boxed{\quad}$$

## ODE solution



## Rigid Body State

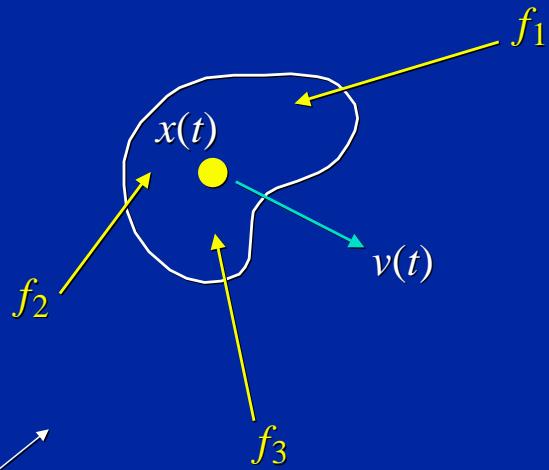


$$\mathbf{X} = \begin{pmatrix} x(t) \\ ? \\ v(t) \\ ? \end{pmatrix}$$

## Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{X} = \frac{d}{dt} \begin{pmatrix} x(t) \\ ? \\ Mv(t) \\ ? \end{pmatrix} = \begin{pmatrix} v(t) \\ ? \\ F(t) \\ ? \end{pmatrix}$$

## Net Force

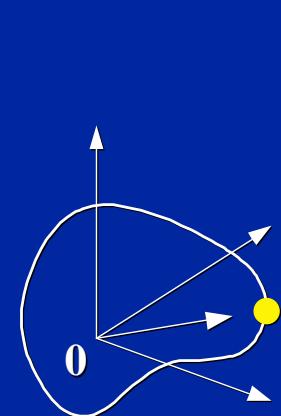


$$F(t) = \sum f_i$$

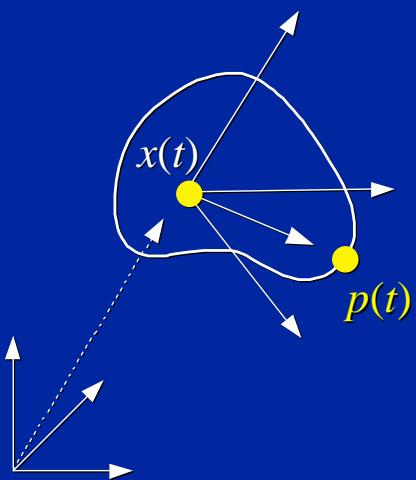
## Orientation

We represent orientation as a rotation matrix  $R(t)$ . Points are transformed from body-space to world-space as:

$$p(t) = R(t)p_0 + x(t)$$



body space



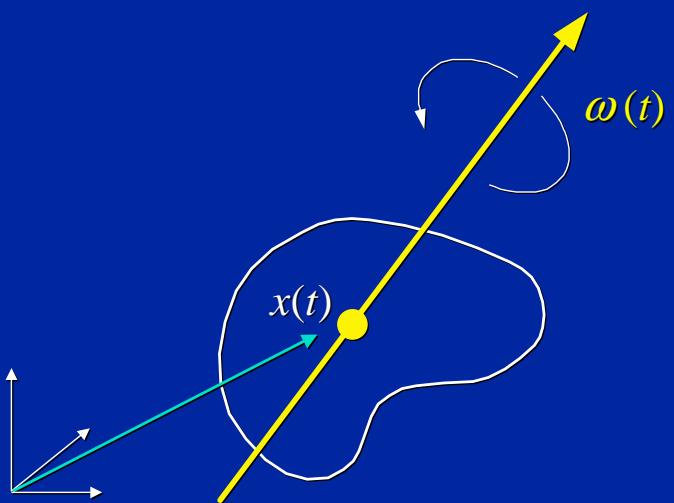
world space

## Angular Velocity

We represent angular velocity as a vector  $\omega(t)$ , which encodes both the axis of the spin and the speed of the spin.

How are  $R(t)$  and  $\omega(t)$  related?

## Angular Velocity Definition



## Angular Velocity

$\dot{R}(t)$  and  $\omega(t)$  are related by

$$\frac{d}{dt} R(t) = \begin{pmatrix} 0 & -\omega_z(t) & \omega_y(t) \\ \omega_z(t) & 0 & -\omega_x(t) \\ -\omega_y(t) & \omega_x(t) & 0 \end{pmatrix} R(t)$$

$\omega(t)^*$  is a shorthand for the above matrix

## Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{X} = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ Mv(t) \\ \boxed{<\omega(t)>} \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ ? \end{pmatrix}$$

Need to relate  $\dot{\omega}(t)$  and mass distribution to  $F(t)$ .

## Inertia Tensor

$$I(t) = \begin{pmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{pmatrix}$$

diagonal terms

$$I_{xx} = M \int_V (y^2 + z^2) dV \quad I_{xy} = -M \int_V xy \, dV$$

off-diagonal terms

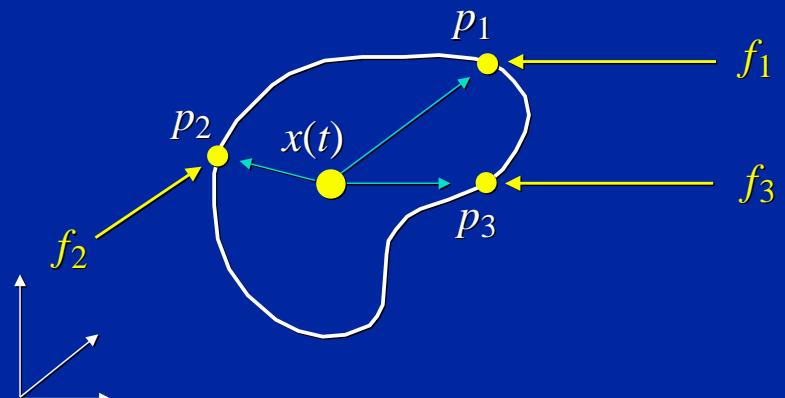
## Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{X} = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ \boxed{Mv(t)} \\ \boxed{I(t)\omega(t)} \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ \tau(t) \end{pmatrix}$$

$P(t)$  – linear momentum

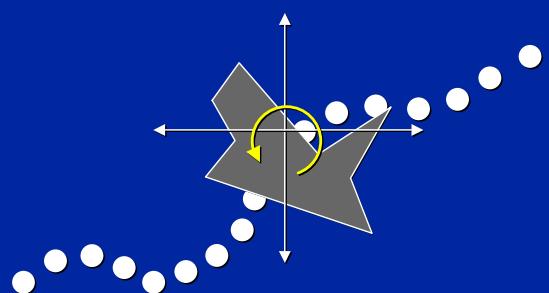
$L(t)$  – angular momentum

## Net Torque



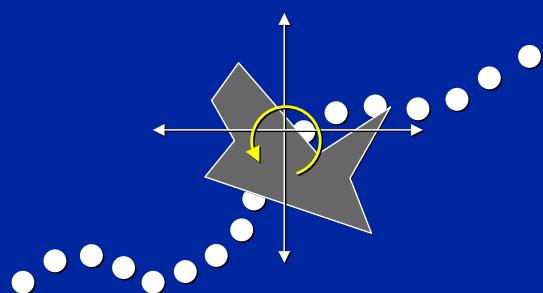
$$\tau(t) = \sum (p_i - x(t)) \times f_i$$

Inertia Tensors Vary in World Space...



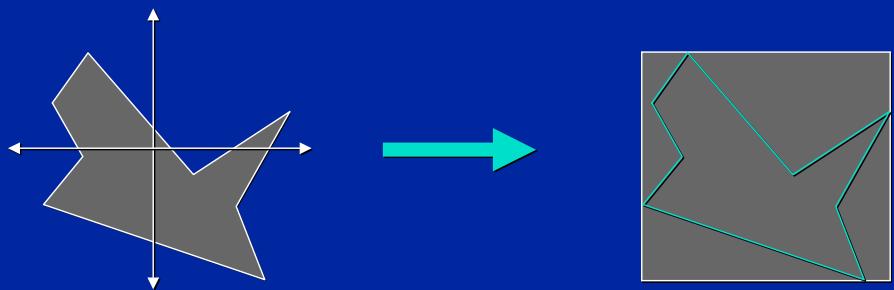
$$I_{xx} = M \int_V (y^2 + z^2) dV \quad I_{xy} = -M \int_V xy \, dV$$

... but are Constant in Body Space



$$I(t) = R(t) I_{\text{body}} R(t)^T$$

## Approximating $I_{\text{body}}$ —Bounding Boxes

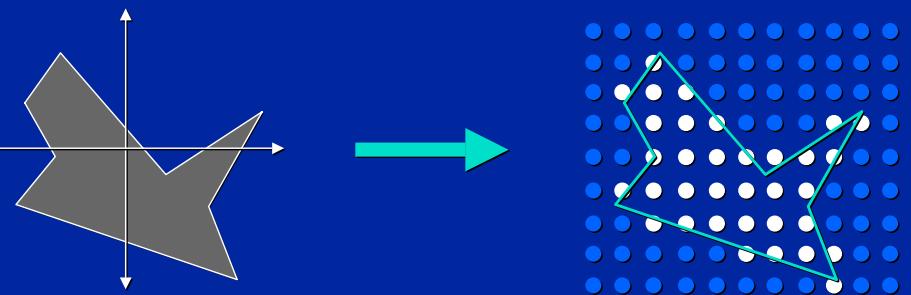


Pros: Simple.

Cons: Bounding box may not be a good fit.

Inaccurate.

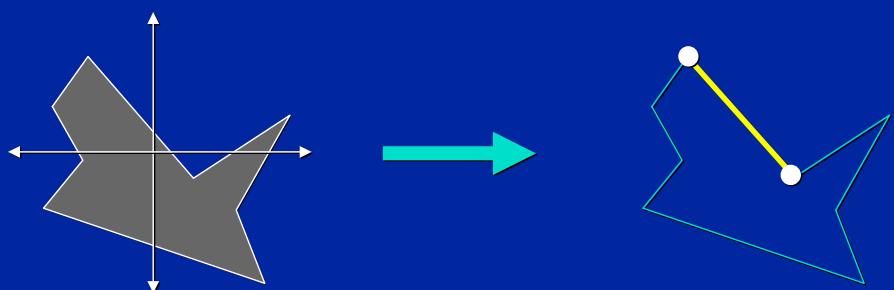
## Approximating $I_{\text{body}}$ —Point Sampling



Pros: Simple, fairly accurate, no B-rep needed.

Cons: Expensive, requires volume test.

## Computing $I_{\text{body}}$ — Green's Theorem (Twice!)



Pros: Simple, exact, no volumes needed.

Cons: Requires B-rep.

Code: <http://www.acm.org/jgt/papers/Mirtich96>

## Rigid Body Equation of Motion

$$\frac{d}{dt} \mathbf{X} = \frac{d}{dt} \begin{pmatrix} \mathbf{x}(t) \\ \mathbf{R}(t) \\ \boxed{\mathbf{Mv}(t)} \\ \boxed{\mathbf{I}(t)\boldsymbol{\omega}(t)} \end{pmatrix} = \begin{pmatrix} \mathbf{v}(t) \\ \boldsymbol{\omega}(t)^* \mathbf{R}(t) \\ \mathbf{F}(t) \\ \boldsymbol{\tau}(t) \end{pmatrix}$$

$\mathbf{P}(t)$  – linear momentum

$\mathbf{L}(t)$  – angular momentum